

# On Poisson–Lie T–plurality of boundary conditions

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# Outline

- 1 Elements of Poisson–Lie T–plurality
- 2 Consistent boundary conditions
- 3 Poisson–Lie T–plurality transformation of the gluing operator

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# Elements of Poisson–Lie T–plurality of $\sigma$ –models

The  $\sigma$ –model given by the **action**

$$S_F[g] = \int_{\Sigma} d^2x \rho_{-}(g) \cdot F(g) \cdot \rho_{+}(g)^t = \int_{\Sigma} d^2x \partial_{-}\phi^{\mu} \mathcal{F}_{\mu\nu}(\phi) \partial_{+}\phi^{\nu} \quad (1)$$

where the map  $g$  maps  $\Sigma = \langle 0, \pi \rangle \times \mathbb{R}$  into the group  $G$  whose Lie algebra has basis  $\{T_a\}$ ,

$$\rho_{\pm}(g)^a \equiv (\partial_{\pm} g g^{-1})^a = \partial_{\pm} \phi^{\mu} e_{\mu}{}^a(g), \quad (\partial_{\pm} g g^{-1}) = \rho_{\pm}(g) \cdot T$$

$\phi^{\mu} : \Sigma \rightarrow \mathbb{R}^{\dim G}$  is the same map as  $g$  but written in some group coordinates, and  $x_{+}, x_{-}$  are the light–cone coordinates on Minkowski  $\mathbb{R}^2$   $\tau = x_{+} + x_{-}$ ,  $\sigma = x_{+} - x_{-}$ .

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$$S_F[g] = \int_{\Sigma} d^2x \rho_-(g) \cdot F(g) \cdot \rho_+(g)^t = \int_{\Sigma} d^2x \partial_- \phi^\mu \mathcal{F}_{\mu\nu}(\phi) \partial_+ \phi^\nu \quad (1)$$

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# The basic idea of Poisson–Lie T–duality

C. Klimčík and P. Ševera, *Phys. Lett. B* 351 (1995) 455.

Under certain conditions the equations of motion in the bulk of the  $\sigma$ –model can be written as equations on

Drinfel'd double

$(G|\tilde{G})$  – Lie group  $D$  whose Lie algebra  $\mathfrak{d}$  admits a decomposition  $\mathfrak{d} = \mathfrak{g} \dot{+} \tilde{\mathfrak{g}}$  into a pair of subalgebras maximally isotropic with respect to a symmetric ad-invariant nondegenerate bilinear form  $\langle \cdot, \cdot \rangle$ .

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If the metric together with the B–field are such that

$$F(g) = (E_0^{-1} + \Pi(g))^{-1}, \quad \Pi(g) = b(g) \cdot a(g)^{-1} = -\Pi(g)^t, \quad (2)$$

then the bulk equations of motion of the  $\sigma$ –model can be formulated as the equations on the Drinfel'd double

$$\langle \partial_{\pm} l l^{-1}, \mathcal{E}^{\pm} \rangle = 0,$$

where  $l = g\tilde{h} \in D$ ,  $g \in G$ ,  $\tilde{h} \in \tilde{G}$  and

$$\mathcal{E}^+ = \text{span} \left( T + E_0 \cdot \tilde{T} \right), \quad \mathcal{E}^- = \text{span} \left( T - E_0^t \cdot \tilde{T} \right)$$

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# Poisson–Lie T–plurality

R. von Unge, J. High En. Phys. 02:07 (2002) 014.

Main idea:

In general there are **several decompositions** (Manin triples) of a Drinfel'd double.

Let  $\hat{\mathfrak{g}} + \bar{\mathfrak{g}}$  be another decomposition of the Lie algebra  $\mathfrak{d}$  into maximal isotropic subalgebras. The dual bases of  $\mathfrak{g}, \tilde{\mathfrak{g}}$  and  $\hat{\mathfrak{g}}, \bar{\mathfrak{g}}$  are related by the linear transformation

$$\begin{pmatrix} T \\ \tilde{T} \end{pmatrix} = \begin{pmatrix} p & q \\ r & s \end{pmatrix} \begin{pmatrix} \hat{T} \\ \bar{T} \end{pmatrix}, \quad (3)$$

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The  $\sigma$ –model related to (1) by the Poisson–Lie T–plurality

is defined analogously but with

$$\begin{aligned}\widehat{F}(\widehat{g}) &= (\widehat{E}_0^{-1} + \widehat{\Pi}(\widehat{g}))^{-1}, & \widehat{\Pi}(\widehat{g}) &= \widehat{b}(\widehat{g}) \cdot \widehat{a}(\widehat{g})^{-1} = -\widehat{\Pi}(\widehat{g})^t, \\ \widehat{E}_0 &= (p + E_0 \cdot r)^{-1} \cdot (q + E_0 \cdot s)\end{aligned}$$

The relation between the classical solutions of equations of motion in the bulk of the two  $\sigma$ –models is obtained from two possible decompositions of  $l \in D$

$$l = g\tilde{h} = \widehat{g}\bar{h}$$

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Does a solution with well–defined boundary conditions transform into another one?

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Does a solution with well–defined boundary conditions transform into another one?**

# Gluing operator

## The gluing operator $\mathcal{R}$

We impose the boundary condition in the form

$$\partial_- g|_{\sigma=0,\pi} = \mathcal{R} \partial_+ g|_{\sigma=0,\pi} \quad (4)$$

Explicitly we write in coordinates or in a frame e.g.

$$\partial_- \phi|_{\sigma=0,\pi} = \partial_+ \phi \cdot R_\phi|_{\sigma=0,\pi}, \quad \rho_-(g)|_{\sigma=0,\pi} = \rho_+(g) \cdot R_\rho|_{\sigma=0,\pi} \quad (5)$$

We define the **Dirichlet projector**  $\mathcal{Q}$  that projects vectors onto the space normal to the D–brane  $\equiv -1$  eigenspace of  $\mathcal{R}$  and **Neumann projector**  $\mathcal{N}$  that projects onto the tangent space of the brane. The corresponding matrices  $Q, N$  are given by

$$Q^2 = Q, \quad Q \cdot R = R \cdot Q = -Q, \quad N = \mathbf{1} - Q. \quad (6)$$

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# Consistency conditions on gluing operator

In addition to (6) we want the following conditions to hold, originally derived in [C. Albertsson, U. Lindström and M. Zabzine, Nucl. Phys. B 678 \(2004\) 295, \[hep-th/0202069\]](#) (in SUSY setting)

- **conformal** – to be consistent with the conformal constraint  $\mathcal{T}_{++}|_{\sigma=0,\pi} = \mathcal{T}_{--}|_{\sigma=0,\pi}$  we need

$$R \cdot (\mathcal{F} + \mathcal{F}^t) \cdot R^t = (\mathcal{F} + \mathcal{F}^t) \quad (7)$$

- **orthogonality** – Neumann and Dirichlet directions must be indeed orthogonal

$$N \cdot (\mathcal{F} + \mathcal{F}^t) \cdot Q^t = 0 \quad (8)$$



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# Consistency conditions on gluing operator

- **integrability** –  $\text{Im}(\mathcal{N})$  must form an integrable distribution, its integral submanifolds being the D–branes

$$N_{\kappa}^{\mu} N_{\lambda}^{\nu} \partial_{[\mu} N_{\nu]}^{\rho} = 0 \quad (9)$$

- **equivalence with the action principle** – the boundary condition should be equivalent to the vanishing variation of the action on the boundary

$$N \cdot (\mathcal{F} - \mathcal{F}^t \cdot R^t) = 0 \quad (10)$$

(is equivalent to the orthogonality condition together with

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## PL T–plurality transformation of the gluing operator

We have found that the transformed solution  $\hat{g}$  satisfies

$$\hat{\rho}_-(\hat{g})|_{\sigma=0,\pi} = \hat{\rho}_+(\hat{g}) \cdot \hat{R}_\rho|_{\sigma=0,\pi} \quad (11)$$

where the transformed gluing operator is

$$\hat{R}_\rho = \hat{F}^t(\hat{g}) \cdot M_-^{-1} \cdot F^{-t}(g) \cdot R_\rho(g) \cdot F(g) \cdot M_+ \cdot \hat{F}^{-1}(\hat{g}), \quad (12)$$

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## Good news

- The transformed gluing operator  $\widehat{R}_\rho$  is found explicitly.
- $\widehat{R}_\rho$  satisfies the conformal condition (7) if and only if the original  $R_\rho$  does (proven)
- $\widehat{R}_\rho$  allows the definition of projectors (6) and satisfies the orthogonality condition (8) if and only if the original  $R_\rho$  does in all the examples investigated for the transitions inside the six–dimensional Drinfel'd doubles ( $Bianchi\ 5 \mid \mathbb{R}^3 \simeq (Bianchi\ 6_0 \mid \mathbb{R}^3)$ ) and the semiabelian four–dimensional Drinfel'd double ( $af(1) \mid \mathbb{R}^2 \simeq (af(1) \mid af(1))$ ).



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## Not so good news

- $\widehat{R}_\rho$  defined by (12) may depend not only on  $\widehat{g}$  but also on  $g$  and consequently on  $\bar{g}$ .

**Solution:**  $\widehat{R}_\rho$  is function of  $\widehat{g}$  only if the matrix–valued function  $C(g) = F^{-t}(g) \cdot R_\rho(g) \cdot F(g)$  extended to a function on the whole Drinfel'd double as  $C(g\tilde{h}) = C(g)$  satisfies

$$C(\widehat{g}\bar{h}) = C(\widehat{g}).$$

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## Bad news

- The integrability condition (9) and equivalence with the action principle (10) are not preserved under the PL T–plurality transformation. (Explicit counterexamples were found.) Sometimes  $\widehat{R}_\rho$  satisfies (9),(10) only in specific points or submanifolds of  $\widehat{G}$ . How to interpret this? We don't know yet.

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# Thank you for you attention

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